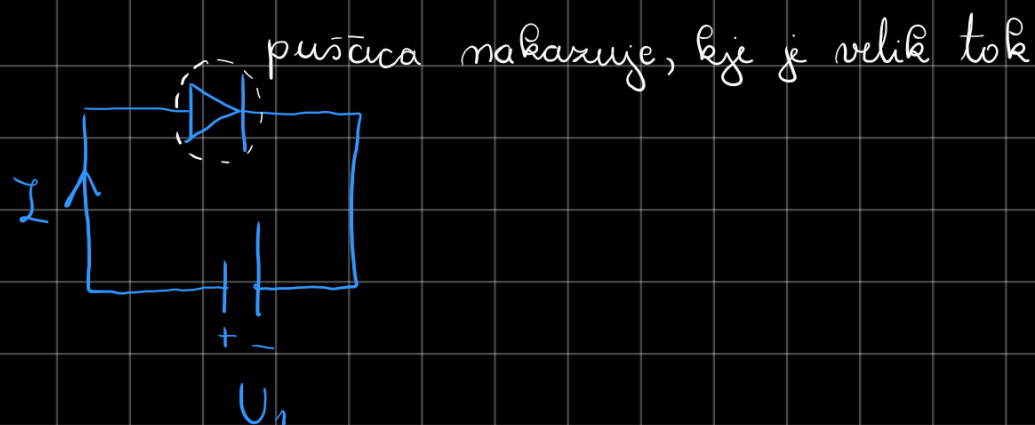


Nadaljevanje od zadnjice

4/18

1) Najprej vzamemo idealno diodo

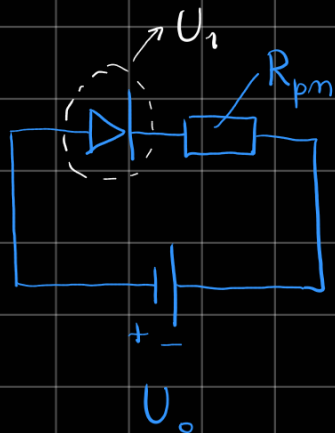


$$I = I_0 (e^{\beta e_0 U} - 1)$$

// območje za U_1

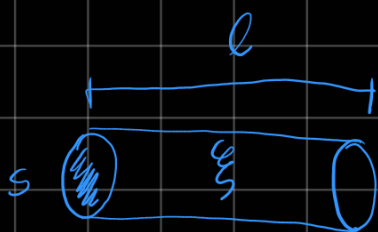
$$U_1 = \frac{1}{\beta e_0} \ln \left[\frac{I}{I_0} + 1 \right] = 0.24V$$

2)



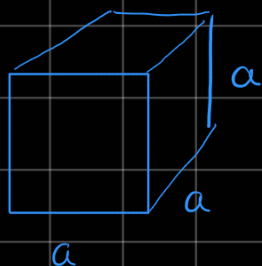
$$U_0 = U_1 + IR$$

// zanima nas R , ker U_0 in U_1 poznamo



$$R = \frac{\xi l}{S}$$

//za kocko je



$$R = \frac{\xi a}{a^2} = \frac{1}{3a}$$

$$R_{pm} = R_p + R_m = \frac{1}{a} \left[\frac{1}{R_p} + \frac{1}{R_m} \right] = 7 \Omega$$

U_0 je napetost vira, ki ga moramo priklopiti

$$U_0 = U_1 + R \cdot I = 0.31V$$

" 0.24V
" 10mA

Jedra

Gustota v jedru je večinoma konst. Vezavna en. je premo sorazmerna s številom gradnikov ($2 \times$ št. gradnikov $\rightarrow 2 \times E_v$)

Velikost je reda $\sim$ fm

jedro: $(p^+ n^0)$

A... oznaka št. gradnikov

↕ mamo št. ($A = Z + N$, N je # neutronov)

Jedro elementa

$\begin{pmatrix} A \\ Z \end{pmatrix} X$

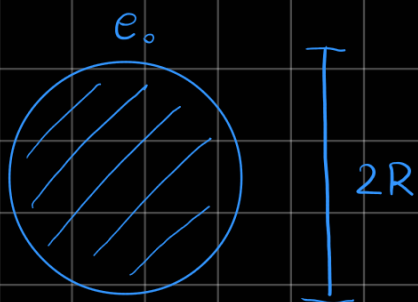
↕ vrstno št. (# protonov)

// predpostavimus, da je jedro iz trdne mase (ni gradienti so tudi
 // empirično lahko iz $\rho = \text{konst.}$ zapišemo radij jedra kroglju)

$$r_j = r_1 A^{\frac{1}{3}}, \quad r_1 = 1.1 \text{ fm}$$

Elektrostatika

Čopček VI/2



// najbolj je enakomerno razmazano

$$E_{el} = ? \quad \phi(r=0) = ?$$

// iz Gaussovega zakona dobimo $\vec{E}$ zunaj in znotraj krogle

$$1) \text{ Gauss} \rightarrow \vec{E}(\vec{r})$$

$$2) \vec{E}(\vec{r}) \rightarrow \phi(\vec{r}) \quad // \text{potencial}$$

$$3) \phi(\vec{r}) \rightarrow E_{el} \quad // \text{elektrostatska en.}$$

$$1) e = \epsilon \oint \vec{E} d\vec{S} \quad d\vec{S} \text{ so sfere (mo v 3D)}$$

$$E(r) = \frac{e(r)}{4\pi r^2 \epsilon_0}$$

$$a) r \leq R$$

$$e = q_e V \quad // \text{u} \text{ } e = e(r) \text{ integral } \int_V e(r) d^3\vec{r}$$

$$e = \frac{\overbrace{e_0}^{g_e}}{\frac{4}{3}\pi R_0^3} \underbrace{\frac{4}{3}\pi r^3}_V$$

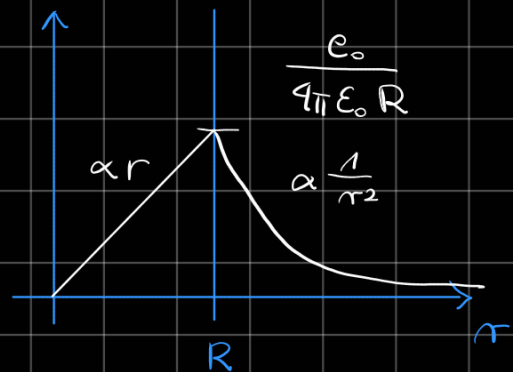
$$E_{<R}(e) = \frac{e_0}{4\pi\epsilon_0 R^3} r$$

b) $r > R$

točkarit nalagaj, $e = \text{konst.}$

$$E(r) = \frac{e(r)}{4\pi\epsilon_0 r^2} \Rightarrow E_{>R}(e) = \frac{e_0}{4\pi\epsilon_0 r^2}$$

$$E(r) = \begin{cases} E_{<R}, & r \leq R \\ E_{>R}, & r > R \end{cases}$$



2) potencial $g(r) = ?$

$$\vec{E}(\vec{r}) = -\nabla g(\vec{r})$$

$$g(\vec{r}) = -\int E(\vec{r}) d\vec{s}$$

// ∇ v kartezičinih $\rightarrow$ // ∇ v sferični

$$\nabla = \left(\frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right) \rightarrow \nabla = \frac{\partial}{\partial r} \hat{e}_r + \frac{1}{r} \frac{\partial}{\partial \theta} \hat{e}_\theta + \frac{1}{r \sin \theta} \frac{\partial}{\partial \phi} \hat{e}_\phi$$

// dila $\hat{e}_\theta$ in $\hat{e}_\varphi$ odpadita, ker $\vec{E} = E(\hat{e}_r)$

a) $r > R$

// nido ločeni int

$$\varphi_{>R}(r) = - \int \frac{e_0}{4\pi\epsilon_0 r^2} dr = \frac{e_0}{4\pi\epsilon_0 r} + A$$

b) $r \leq R$

$$\varphi_{<R}(r) = - \int \frac{e_0 r}{4\pi\epsilon_0 R^3} dr = - \frac{e_0 r^2}{8\pi\epsilon_0 R^3} + B$$

a. 1) $\varphi(r \rightarrow \infty) \rightarrow 0 \Rightarrow A = 0$

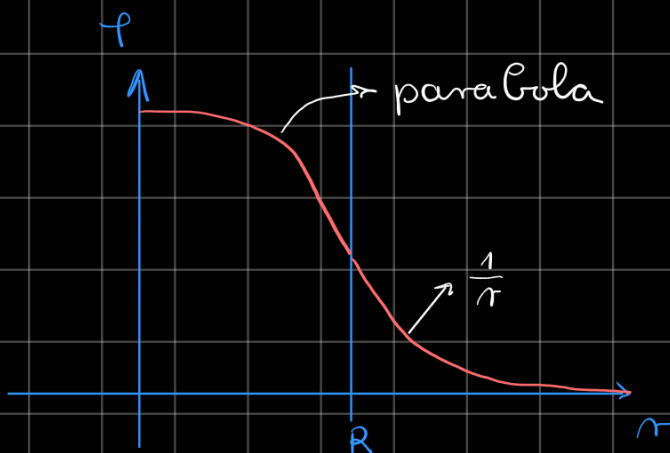
b. 1) $\varphi_{<R}(r=R) = \varphi_{>R}(r=R)$

$$- \frac{e_0 r^2}{8\pi\epsilon_0 R} + B = \frac{e_0}{4\pi\epsilon_0 R}$$

$$\Rightarrow B = \frac{3e_0}{8\pi\epsilon_0 R}$$

$$\varphi(r=0) = \frac{3e_0}{8\pi\epsilon_0 R}$$

$$\varphi_{>R}(r) = \frac{e_0}{8\pi\epsilon_0 R} \left[3 - \left(\frac{r}{R} \right)^2 \right]$$



3) Elektrostatika energija

$$\begin{array}{c} \text{-----} \quad \rho(\vec{r}) \\ \text{-----} \\ \text{točkasti naboj} \\ \bullet \vec{r} = \vec{r}_0 \\ \text{-----} \\ \text{-----} \end{array}$$

U mišamo točkastega naboja, ampak nitočkastega, razbijemo na infinitesimalne dele

$$E_{el} = e \rho(\vec{r} = \vec{r}_0)$$

// za jedro $E_{el} = \left(\frac{1}{2} \right) \int_V \rho(\vec{r}) d\vec{r}$

parska interakcija, deluje in izračunamo je na par, mas pa xanima samo za krogljico

$$e(r) = \frac{e_0 r^3}{R^3} \Rightarrow \frac{de}{dr} = \frac{3e_0 r^2}{R^3}$$

$$\begin{aligned} E_{el} &= \frac{1}{2} \frac{e_0}{8\pi\epsilon_0 R} \frac{3e_0}{R^3} \int_0^R dr r^2 \left[3 - \left(\frac{r}{R} \right)^2 \right] \\ &= \frac{3e^2}{16\pi\epsilon_0 R^4} \left[\left. 3 \frac{x^3}{3} \right|_0^1 - \left. \frac{x^5}{5} \right|_0^1 \right] \end{aligned}$$

$$\begin{aligned} x &= \frac{r}{R} \\ \frac{dx}{dr} &= \frac{1}{R} \end{aligned}$$

$$E_d = \frac{3e^2}{20\pi\epsilon_0 R}$$

// za običtek $R \sim \text{fm} \Rightarrow E_d \sim \text{MeV}$

$$E_d = \frac{3}{5} \frac{e^2}{4\pi\epsilon_0 \hbar c} \frac{\hbar c}{R}$$

α

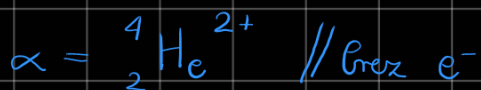
$\xrightarrow{197 \text{ MeV fm}}$
 $\sim \text{fm}$

$\alpha \doteq \frac{1}{137}$ // konstanta fine
strukture

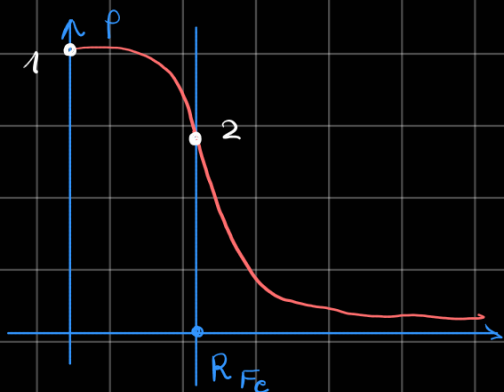
VI / 3



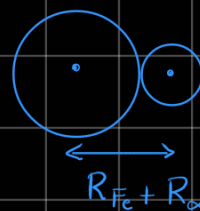
// ločimo kinetično en. α delca



Skica

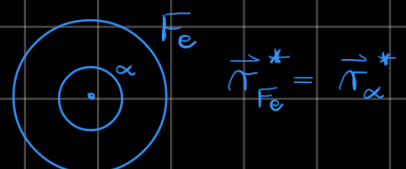


// pri 2 α jedri dotikata



// tukaj bomo
vzeli potencial, ki
je zunaj jedra

// pri 1 α mediji prenikata



Fe:

$$A = 56, \quad Z = 26$$

α :

$$A = 4, \quad Z = 2$$

// empirična formula od prej $r_j = r_1 A^{\frac{1}{3}}$

$$R_{Fe} = 4'21 \text{ fm}$$

$$R_{\alpha} = 1'75 \text{ fm}$$

1) $E_1 = P(r=0) \cdot e_{\alpha}$ // predpostavim, da je α točkasto

$$= \frac{3 Z_{Fe} Z_{\alpha} e^2}{8\pi \epsilon_0 R_{Fe}} = \frac{3}{2} Z_{Fe} Z_{\alpha} \alpha \cdot \frac{\hbar c}{R_{Fe}} = 26'6 \text{ MeV}$$

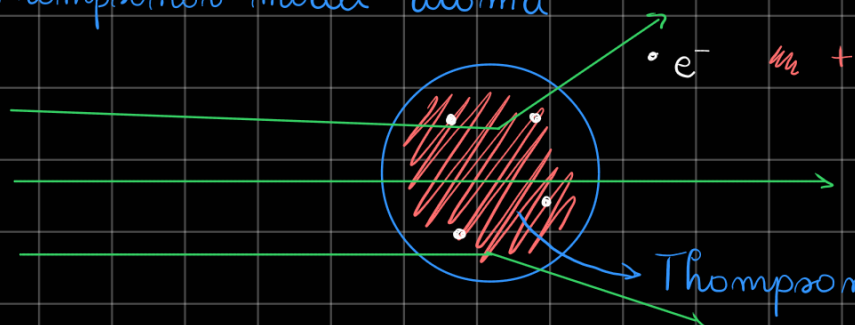
$$2) E_2 = \frac{Z_{Fe} Z_{\alpha} e^2}{4\pi \epsilon_0 (R_{Fe} + R_{\alpha})} = 12'5 \text{ MeV}$$

Rutherfordovo sipanje

Pred Rutherfordom je bil Thompsonov puding (homogena masa s točkastimi naboji)

Po poskusu: jedro je majhen del atoma

Thompsonov model atoma



// maksimalni kot sipanja
 $\theta = 20^\circ$

Thompsonov puding

